# Close to Close Packing 

Oliver Penrose ${ }^{1}$ and George Stell ${ }^{2}$

Received September 20, 1999


#### Abstract

For various lattice gas models with nearest neighbor exclusion (and, in one case, second-nearest neighbor exclusion as well), we obtain lower bounds on $m$, the average number of particles on the nonexcluded lattice sites closest to a given particle. They are all of the form $m / m_{c p} \geqslant 1-$ const $\cdot\left(N_{c p} / N-1\right)$, where $N$ is the number of occupied sites, $m_{c p}$ is the value of $m$ at close packing, and $N_{c p}$ is the value of $N$ at close packing. An analogous result exists for hard disks in the plane.


KEY WORDS: Close packing; lattice models; inequalities; nearest neighbor exclusion; hard disks.

## 1. INTRODUCTION

In a classic piece of work, ${ }^{(3)}$ L. Fejes Tóth gave a proof that the highestdensity packing of non-overlapping disks in a plane is, as one expects intuitively, the hexagonal close packing. His ideas were used by the present authors ${ }^{(2)}$ to obtain information about the packing of such disks at densities slightly below the close-packing density. Defining $m(r)$ to be the average number of disks whose centres lie within a distance $r$ of a given disk, we showed that $m(r)$ satisfies an inequality

$$
\begin{equation*}
1 \geqslant \frac{m(r)}{6} \geqslant 1-\frac{A / A_{c p}-1}{\left(r^{2} / a^{2}\right)-1}-\varepsilon \quad \text { if } \quad 1<r / a<\frac{1}{2} \operatorname{cosec} \frac{\pi}{7}=1.15 \ldots \tag{1}
\end{equation*}
$$

Here $a$ is the diameter of the disks, $A$ is the area of the (hexagonal) region available to the disks, $A_{c p}=\frac{1}{2} \sqrt{3} N a^{2} \approx 0.866 N a^{2}$ is the value of $A$ at hexagonal

[^0]close packing, $N$ is the number of disks, and $\varepsilon$ denotes a correction term which is much smaller than the term shown if $A / a^{2} \gg 1$ and $r / a-1 \ll 1$.

This result can be used, for example, to obtain upper and lower bounds on the energy of a system of hard disks with a sufficiently shortrange square-shoulder or square-well interaction

$$
U(r)= \begin{cases}+\infty & (0 \leqslant r<a)  \tag{2}\\ U_{0} & (a \leqslant r<b) \\ 0 & (b \leqslant r)\end{cases}
$$

where $U_{0}$ and $b$ are constants, with $a<b<\frac{1}{2} a \operatorname{cosec}(\pi / 7)$. These bounds provide an upper bound on the error in the formula for the mean energy per particle given by thermodynamic perturbation theory; this upper bound is proportional to the deviation of the density from the close-packing density.

In view of the recent mathematical progress ${ }^{(1)}$ on the three-dimensional analogue of Fejes Tóth's problem (often called Kepler's conjecture), one may hope that a three-dimensional analogue of the inequality (1) may some day be proven. In the present note, however, our ambitions are much more modest: to prove some results analogous to (1) for lattice gases with nearest neighbor exclusion, both in two and three dimensions. We shall consider (though not in that order) a square lattice with nearest neighbor exclusion, a triangular lattice with nearest neighbor exclusion, a cubic lattice with nearest neighbor exclusion, and a cubic lattice with both nearest- and second-nearest neighbor exclusion. In all the cases considered the result we obtain can be written in the form

$$
\begin{equation*}
1 \geqslant \frac{m}{m_{c p}} \geqslant 1-\text { const } \cdot\left(\frac{N_{c p}}{N}-1\right) \tag{3}
\end{equation*}
$$

where $N$ is the number of particles (i.e., occupied lattice sites), $m$ is the average number of particles at the nonexcluded lattice sites nearest to a given particle, and $N_{c p}, m_{c p}$ are the values of $N, m$ at close packing. Except in the case of the triangular lattice, the constant is 1 . Since the fractions $A / A_{c p}$ in (1) and $N_{c p} / N$ in (3) are equal in the thermodynamic limit, there is a close analogy between (1) and (3).

## 2. PLANE SQUARE LATTICE

Consider a plane square lattice with $L$ sites, $N$ of which are occupied, subject to nearest neighbor exclusion. We denote by $m$ the average number of second-nearest neighbors per particle (the number of nearest neighbors
is, of course, zero). To avoid edge effects we take the underlying space to be a torus and give the lattice an even number of sites along each axis, so that a perfect close-packing arrangement is possible.

Let $p_{0}, p_{1}, p_{2}$ denote, respectively, the number of plaquettes with 0,1 , and 2 corners occupied; because of the nearest neighbor exclusion rule, $p_{2}$ is the number of plaquettes with two opposite corners occupied and the other two unoccupied. The number of second-nearest neighbor pairs is then

$$
\begin{equation*}
p_{2}=\frac{1}{2} m N \tag{4}
\end{equation*}
$$

by the definition of $m$.
Since, on this lattice, the total number of plaquettes is equal to the number of sites we have

$$
\begin{equation*}
p_{0}+p_{1}+p_{2}=L \tag{5}
\end{equation*}
$$

Also, since each site meets four plaquettes, the total number of occupied corners of plaquettes is equal to four times the number of occupied sites, and hence

$$
\begin{equation*}
p_{1}+2 p_{2}=4 N \tag{6}
\end{equation*}
$$

At close packing, all the plaquettes have two occupied corners, so that $p_{0}=p_{1}=0, p_{2}=L$. From the above equations, the values of $N$ and $m$ at close packing are

$$
\begin{align*}
& N_{c p}=\frac{1}{2} L  \tag{7}\\
& m_{c p}=4 \tag{8}
\end{align*}
$$

Evidently $m_{c p}$ is the coordination number of the occupied sublattice. From (4) and (6) we find

$$
\begin{equation*}
(4-m) N=p_{1} \tag{9}
\end{equation*}
$$

and, from (7), (5) and (6),

$$
\begin{equation*}
N_{c p}-N=\frac{1}{2} L-N=\frac{1}{2} p_{0}+\frac{1}{4} p_{1} \tag{10}
\end{equation*}
$$

so that

$$
\begin{equation*}
0 \leqslant(4-m) N \leqslant 4\left(N_{c p}-N\right) \tag{11}
\end{equation*}
$$

This gives, as our analogue of (1) for the square lattice with nearest neighbor exclusion,

$$
\begin{equation*}
1 \geqslant \frac{m}{4} \geqslant 1-\left(N_{c p} / N-1\right) \tag{12}
\end{equation*}
$$

## 3. SIMPLE CUBIC LATTICE, NEAREST NEIGHBOR EXCLUSION

Now consider a simple cubic lattice with $L$ sites, $N$ of which are occupied, subject to nearest neighbor exclusion. As before, we denote by $m$ the average number of second-nearest neighbors per particle and take the underlying space to be a (three-dimensional) torus. In the close-packed arrangement, the occupied sites form a face-centred cubic lattice.

As before, let $p_{0}, p_{1}, p_{2}$ denote, respectively, the number of plaquettes with 0,1 , and 2 corners occupied; because of the nearest neighbor exclusion rule, $p_{2}$ is the number of plaquettes with two opposite corners occupied and the other two unoccupied. The number of second-nearest neighbor pairs is again given by (4), but for this lattice, the total number of plaquettes is three times the number of sites so that

$$
p_{0}+p_{1}+p_{2}=3 L
$$

Also, since each site meets 12 plaquettes, the analogue of (6) is

$$
p_{1}+2 p_{2}=12 N
$$

At close packing, all the plaquettes have two occupied corners, so that $p_{0}=p_{1}=0, p_{2}=L$. The values of $N$ and $m$ at close packing are

$$
\begin{align*}
& N_{c p}=\frac{1}{2} L  \tag{15}\\
& m_{c p}=12 \tag{16}
\end{align*}
$$

From (4) and (14) we find

$$
\begin{equation*}
(12-m) N=p_{1} \tag{17}
\end{equation*}
$$

and from (15), (13) and (14),

$$
\begin{equation*}
N_{c p}-N=\frac{1}{2} L-N=\frac{1}{6} p_{0}+\frac{1}{12} p_{1} \tag{18}
\end{equation*}
$$

so that

$$
\begin{equation*}
0 \leqslant(12-m) N \leqslant 12\left(N_{c p}-N\right) \tag{1}
\end{equation*}
$$

This gives, as our analogue of (1) for the cubic lattice with nearest neighbor exclusion,

$$
\begin{equation*}
1 \geqslant \frac{m}{12} \geqslant 1-\left(N_{c p} / N-1\right) \tag{20}
\end{equation*}
$$

## 4. SIMPLE CUBIC LATTICE, NEAREST- AND SECOND-NEAREST NEIGHBOR EFXCLUSION

A similar result also exists for the same lattice but with exclusion on both nearest and second-nearest neighbor sites. In this case the occupied sites in the closest-packed configuration form a body-centred cubic lattice.

This time, instead of plaquettes, we consider the small cubes whose corners are nearest neighbor lattice sites. Let $c_{0}, c_{1}, c_{2}$ denote, respectively, the number of such cubes with 0,1 , and 2 corners occupied; because of the exclusion rule, $c_{2}$ is exactly the number of small cubes with two opposite corners occupied and the other two unoccupied. Let $m^{\prime}$ denote the average number of second-nearest neighbors of a given occupied site, so that the number of second-nearest neighbor pairs is given by

$$
\begin{equation*}
c_{2}=\frac{1}{2} m^{\prime} N \tag{21}
\end{equation*}
$$

Since the total number of small cubes is equal to the number of sites we have

$$
\begin{equation*}
c_{0}+c_{1}+c_{2}=L \tag{22}
\end{equation*}
$$

Also, since each site meets 8 small cubes, the analogue of (6) is now

$$
\begin{equation*}
c_{1}+2 c_{2}=8 N \tag{23}
\end{equation*}
$$

At close packing, all the small cubes have two occupied corners, so that $c_{0}=c_{1}=0, c_{2}=L$. From the above equations, the values of $N$ and $m$ at close packing are

$$
\begin{align*}
& N_{c p}=\frac{1}{4} L  \tag{24}\\
& m_{c p}^{\prime}=8 \tag{25}
\end{align*}
$$

From (21) and (23) we find

$$
\begin{equation*}
\left(8-m^{\prime}\right) N=c_{1} \tag{26}
\end{equation*}
$$

and from (24), (22) and (23),

$$
\begin{equation*}
N_{c p}-N=\frac{1}{4} L-N=\frac{1}{4} c_{0}+\frac{1}{8} c_{1} \tag{27}
\end{equation*}
$$

so that

$$
\begin{equation*}
0 \leqslant\left(8-m^{\prime}\right) N \leqslant 8\left(N_{c p}-N\right) \tag{28}
\end{equation*}
$$

This gives, as our analogue of (1) for the cubic lattice with nearest- and second-nearest neighbor exclusion,

$$
\begin{equation*}
1 \geqslant \frac{m^{\prime}}{8} \geqslant 1-\left(N_{c p} / N-1\right) \tag{29}
\end{equation*}
$$

## 5. TRIANGULAR LATTICE

To obtain an inequality for the triangular lattice with nearest neighbor exclusion, a slightly more complicated argument is necessary. This time the plaquettes are triangles. Using the same notation $p_{0}, p_{1}, p_{2}$ as in Section 2, we see that $p_{2}=0$ because of the exclusion rule, so that the analogues of (5), (7) and (6) are

$$
\begin{align*}
p_{0}+p_{1} & =2 L=6 N_{c p}  \tag{30}\\
p_{1} & =6 \mathrm{~N} \tag{31}
\end{align*}
$$

from which it follows that

$$
\begin{equation*}
p_{0}=6\left(N_{c p}-N\right) \tag{32}
\end{equation*}
$$

Let us define also $s_{0}, s_{1}, s_{2}$ as the number of second-nearest neighbor pairs of sites with, respectively, neither site occupied, one site occupied, or both occupied. They too satisfy relations analogous to (4), (5) and (6); the ones we shall need are

$$
\begin{align*}
s_{1}+2 s_{2} & =6 N  \tag{33}\\
s_{2} & =\frac{1}{2} m N \tag{34}
\end{align*}
$$

from which it follows that

$$
\begin{equation*}
s_{1}=(6-m) N \tag{35}
\end{equation*}
$$

For every second-nearest pair of sites, one of which is occupied and the other occupied, there is one empty plaquette and one singly occupied
plaquette. Since each empty plaquette can belong in this way to at most three second-nearest pairs of this type, we have

$$
\begin{equation*}
s_{1} \leqslant 3 p_{0} \tag{36}
\end{equation*}
$$

From (31), (35) and (36) we find

$$
\begin{equation*}
(6-m) N \leqslant 18\left(N_{c p}-N\right) \tag{37}
\end{equation*}
$$

so that the analogue of (1) for the triangular lattice with nearest neighbor exclusion is (since the $m_{c p}$ for this lattice is 6 )

$$
\begin{equation*}
1 \geqslant \frac{m}{6} \geqslant 1-3\left(N_{c p} / N-1\right) \tag{38}
\end{equation*}
$$

## ACKNOWLEDGMENTS

O.P. thanks DIMACS at Rutgers University for financial support which helped with the writing of this paper. DIMACS is an NSF Science and Technology Center, funded under contract STC-91-1999; it also receives support from the New Jersey Commission on Science and Technology. G.S. acknowledges the support of the Division of Chemical Sciences, Office of Basic Energy Sciences, Office of Science, U.S. Department of Energy.

## REFERENCES

1. Thomas C. Hales, http://www.math.lsa.umich.edu/hales/countdown/, cited by Keith Devlin, http://www.maa.org/devlin/devlin_9_98.html.
2. G. Stell and O. Penrose, Phys. Rev. Lett. 51:1397-1400 (1983). First-order thermodynamic perturbation theory is exact in the two-dimensional close-packing limit.
3. L. Fejes Tóth, Regular Figures (Macmillan, New York, 1964), p. 165.

[^0]:    We dedicate this work to the memory of Ann Stell, beloved wife of GS and beloved friend of OP.
    ${ }^{1}$ Mathematics, Heriot-Watt University, Edinburgh EH14 4AS, United Kingdom.
    ${ }^{2}$ Chemistry, SUNY, Stony Brook, New York.

